

Gravitational Many-Body Problem

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Talk structure

1. Overview of Gravitational Many-Body Problem
 - collisionless and collisional systems
 - Evolution of collisional systems
2. Some recent topics
 - Black hole formation in star clusters
3. Simulation methods

Gravitational Many-Body Problem

- Definition
- Astronomical examples
- Large- N limit — “collisionless” systems
- Two-body (collisional) relaxation
- Evolution of collisional systems

Definition

$$m_i \frac{d^2 x_i}{dt^2} = \sum_{j \neq i} f_{ij} \quad (1)$$

x_i, m_i : position and mass of particle i

f_{ij} : force from particle j to particle i

Gravitational (Newtonian) force:

$$f_{ij} = G m_i m_j \frac{x_j - x_i}{|x_j - x_i|^3}, \quad (2)$$

G : Gravitational constant

Astronomical Examples

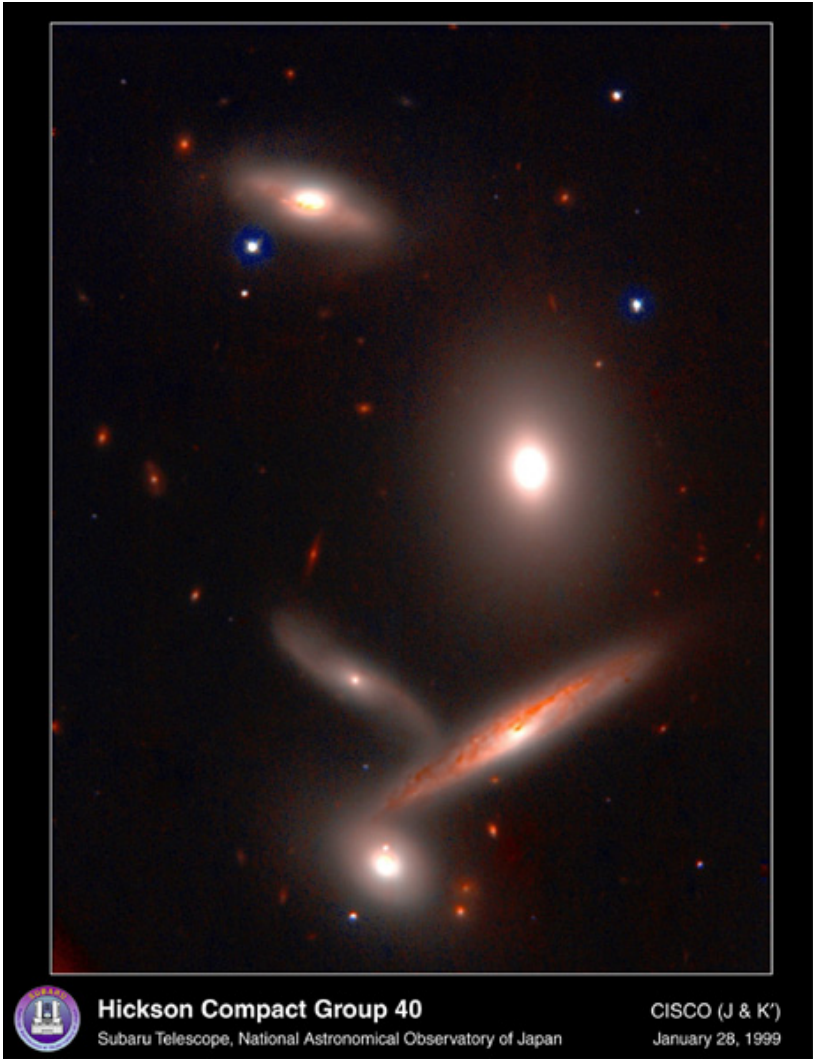
Star Clusters



Galaxy



Galaxy Group



Clusters of Galaxies

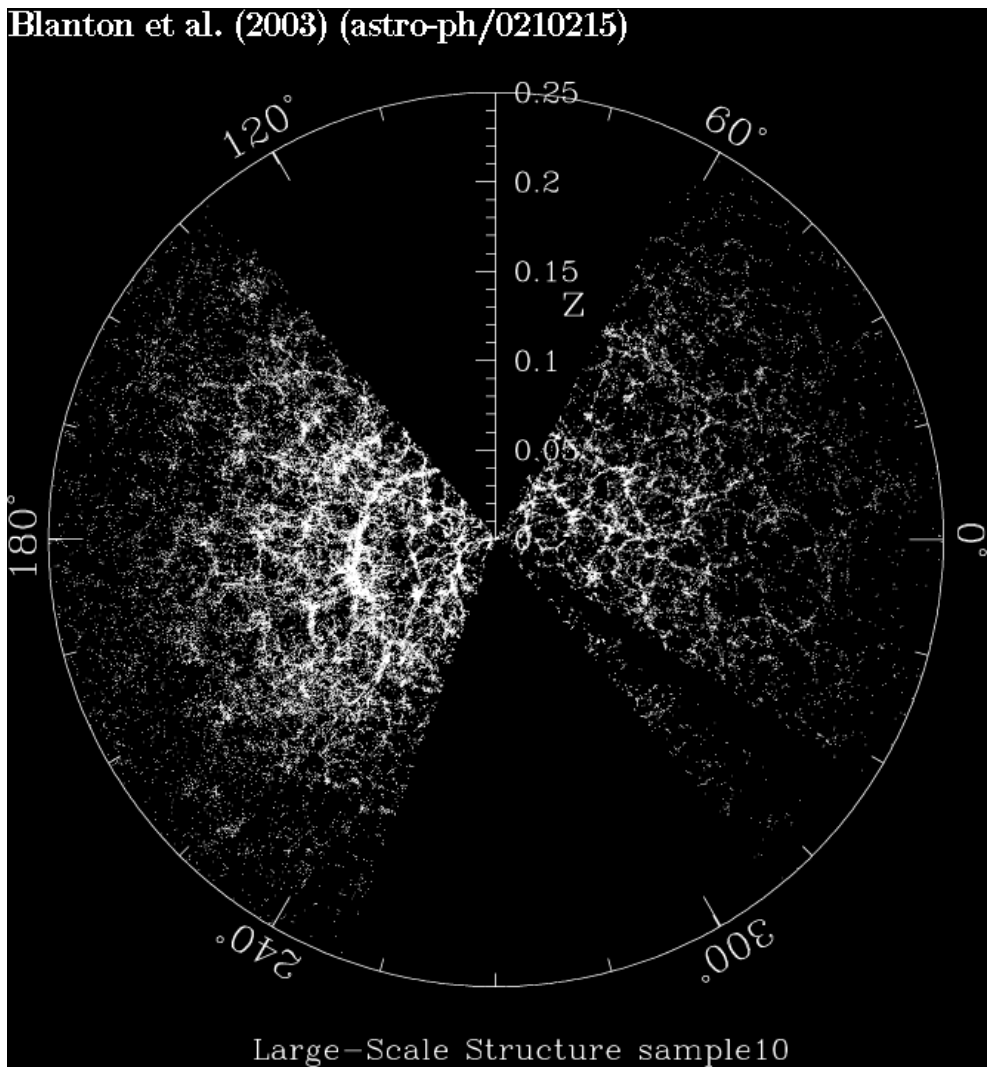


Galaxy Cluster
0024+1645

Yellow: Galaxies
in the cluster

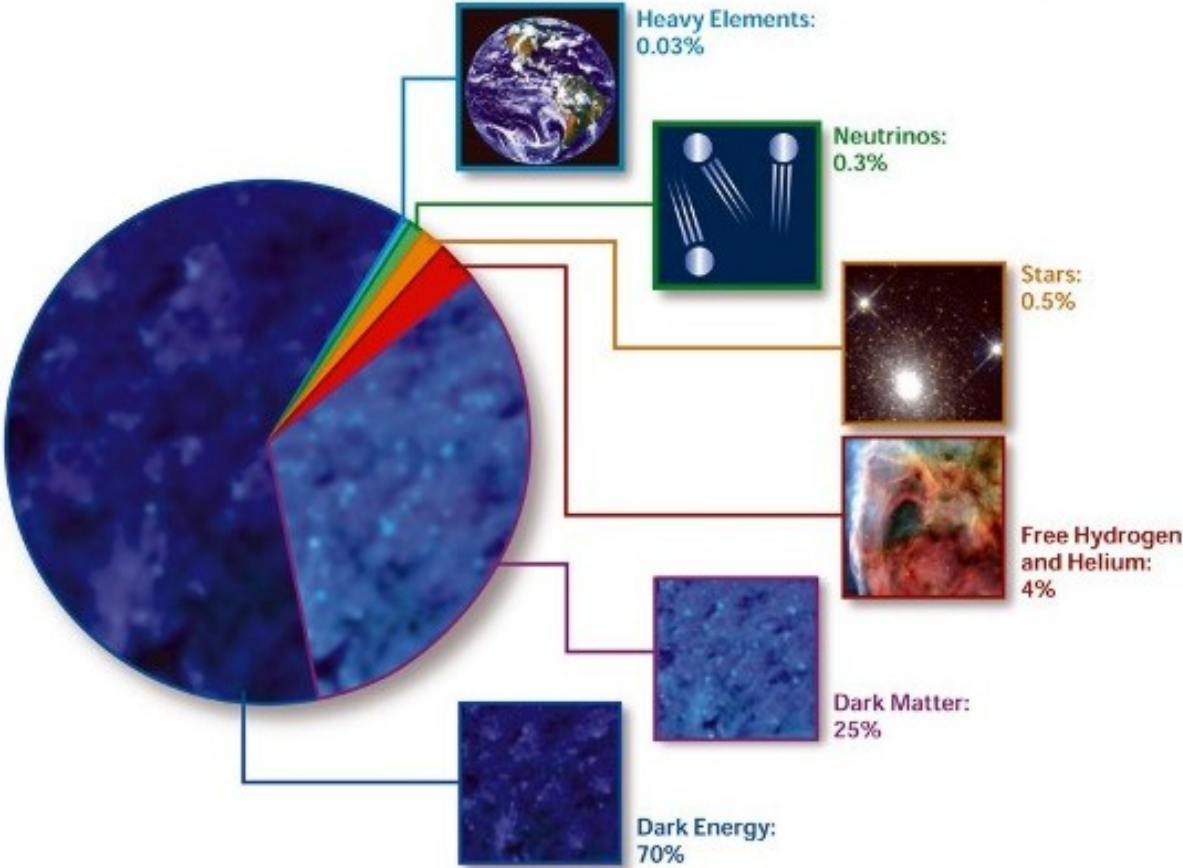
Blue: background
galaxies (enlarged
and stretched
through
gravitational lens
effect)

Large Scale Structure



SDSS Survey
Radius: about
1Gpc (1/3 of the
Universe)

Constituent of the Universe



70%: Interacts with nothing

25%: Interacts only through gravity

Behavior of Dark Matter

Initial conditions

- Power spectrum of density fluctuation
- Cosmological parameters

Fairly well understood (if our assumption on the nature of DM is correct)

We can simulate the structure formation through dark matter dynamics

(Animation)

Large- N Limit

— collisionless systems

Collisionless Boltzmann Equation (CBE)

$$\frac{\partial f}{\partial t} + \boldsymbol{v} \cdot \nabla f - \nabla \Phi \cdot \frac{\partial f}{\partial \boldsymbol{v}} = 0, \quad (3)$$

f : Distribution function in 6-dimensional phase space

Φ : Gravitational potential, given by:

$$\nabla^2 \Phi = 4\pi G \rho. \quad (4)$$

ρ : mass density

$$\rho = m \int d\boldsymbol{v} f, \quad (5)$$

Dynamical Equilibrium

Dynamical Equilibrium: Stationary solution of
CBE+Poisson Eq.

Star clusters, Galaxies: Roughly in DE

Large Scale Structure: NOT in DE
(Dynamically too young)

Jeans' Theorem

Stationary solution of CBE in a given potential ϕ can be expressed as

$$f(\mathbf{x}, \mathbf{v}) = f[I_1(\mathbf{x}, \mathbf{v}), I_2(\mathbf{x}, \mathbf{v}), \dots], \quad (6)$$

where I_1, I_2, \dots are integrals of the motion.

Example: Spherically symmetric system:

$f = f(E, J)$. (energy and total angular momentum)

The distribution function f for a stellar system with finite mass cannot be Maxwellian, because f must be zero for $E > 0$.

Dynamical and Thermal Equilibrium

Dynamical Equilibrium: Stationary state of CBE

No entropy production (no collision term)

→ Not in thermal equilibrium

Not in LOCAL thermal equilibrium either
(Jeans' Theorem)

Collision term

Two-body relaxation

Close encounter between two point-mass particles

Orbit changes: similar to physical collision

Relaxation timescale:

$$t_r \sim \frac{v^3}{Gnm^2 \log \Lambda} \quad (7)$$

$\log \Lambda$: Coulomb Logarithm, $\log \Lambda \sim \log N$

$$t_r \sim \frac{N}{\log N} t_d \quad (8)$$

t_d : Dynamical timescale

Relaxation timescale \propto Number of particles

Thermal (two-body) relaxation

- Important for
 - Small- N systems
 - Systems with small dynamical timescale
- Examples
 - Star clusters
 - Central region of galaxies
 - Clusters of Galaxies
 - Few-body systems (triple etc)

Virial Theorem and negative specific heat

For a stellar system in dynamical equilibrium

$$E = -K = V/2 \quad (9)$$

E : Total energy ($= K + V$)

K : Kinetic energy

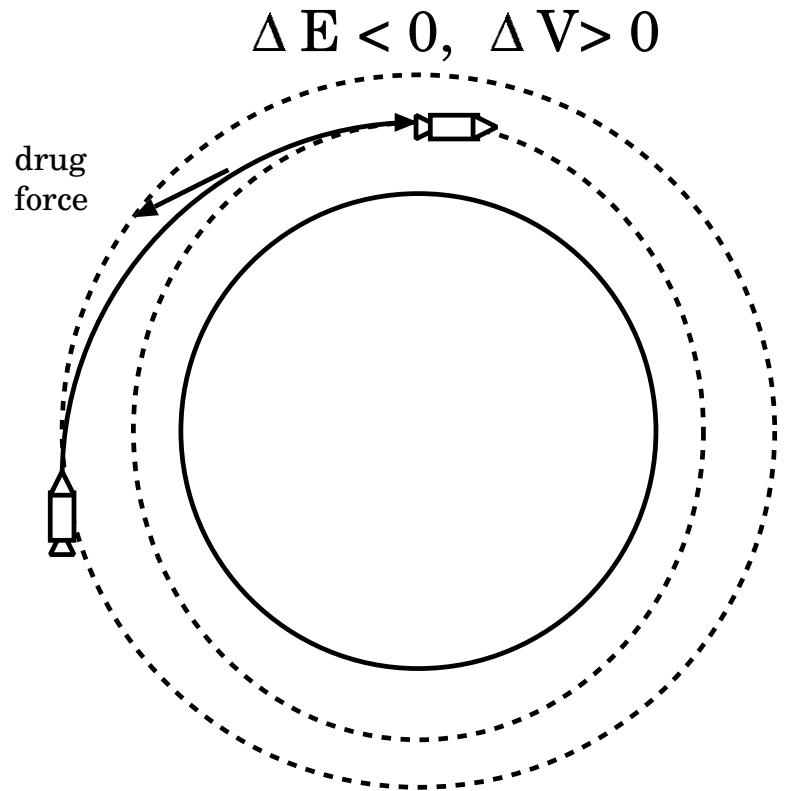
V : Potential energy

If a system loses energy

($\Delta E < 0$),

it becomes hotter

($\Delta K > 0$)



Gravothermal Instability

Isothermal gas in
a spherical adiabatic wall

Radius: R

Mass: M

Central density: ρ_0

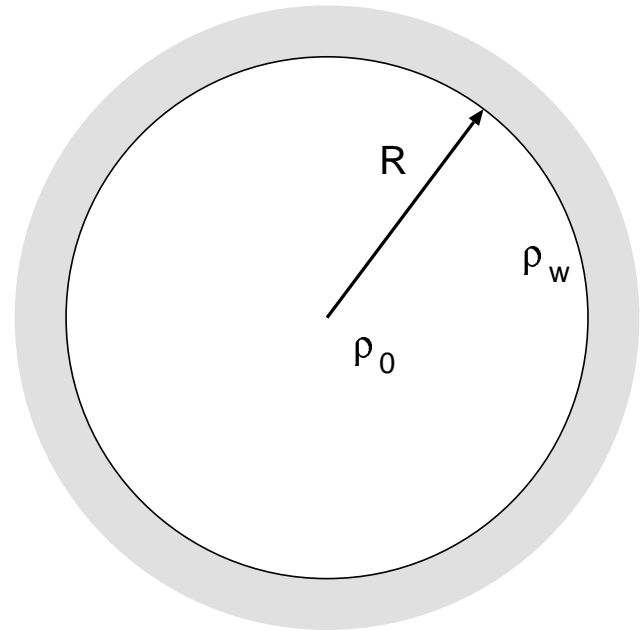
Density at the wall: ρ_w

$$D = \rho_0 / \rho_w$$

$D = 1$: No gravity

$D = \infty$: Singular solution

$$\rho \propto r^{-2}$$



Gravothermal Instability(cont'd)

For large D ($D > 709$):

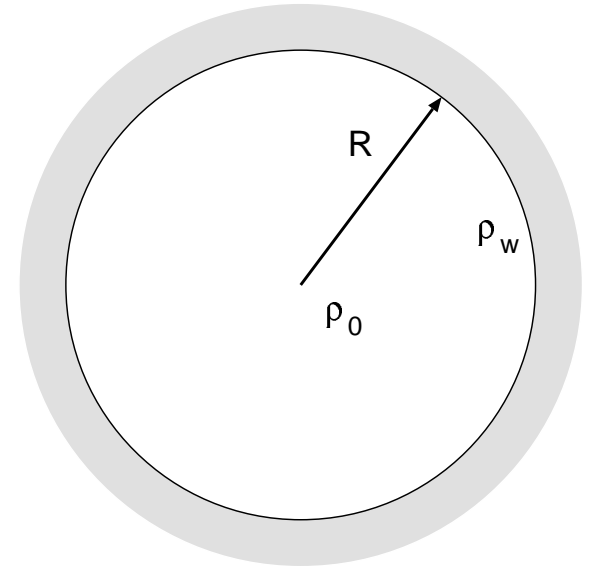
move heat from the central region to
outer Isothermal gas



central region becomes hotter due to
negative specific heat



heat flow grows



Evolution in finite amplitude

Numerical calculation

Hachisu *et al.* (1978) :

Gas dynamics

Lynden-Bell & Eggleton

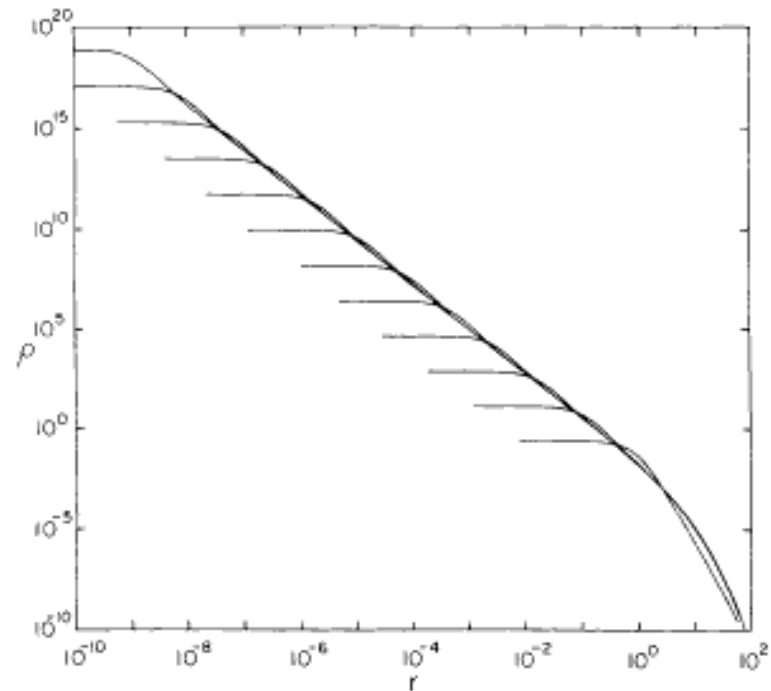
(1980): Self-similar solution

Cohn (1980): Direct

integration of

orbit-averaged

Fokker-Planck equation



Energy Production

Gravitational contraction of star:
halted by nuclear fusion

Stellar system: Three-body binary formation

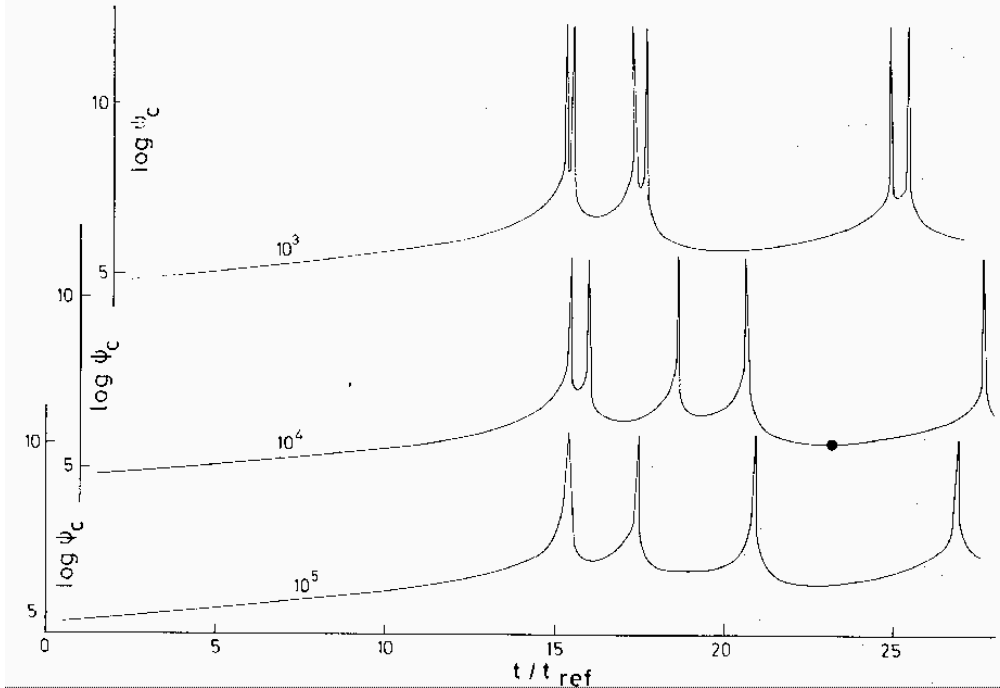
Close encounter of three particles \rightarrow
One binary + one single star:
Statistically most likely outcome

Animation

Three particles

Gravothermal Oscillation

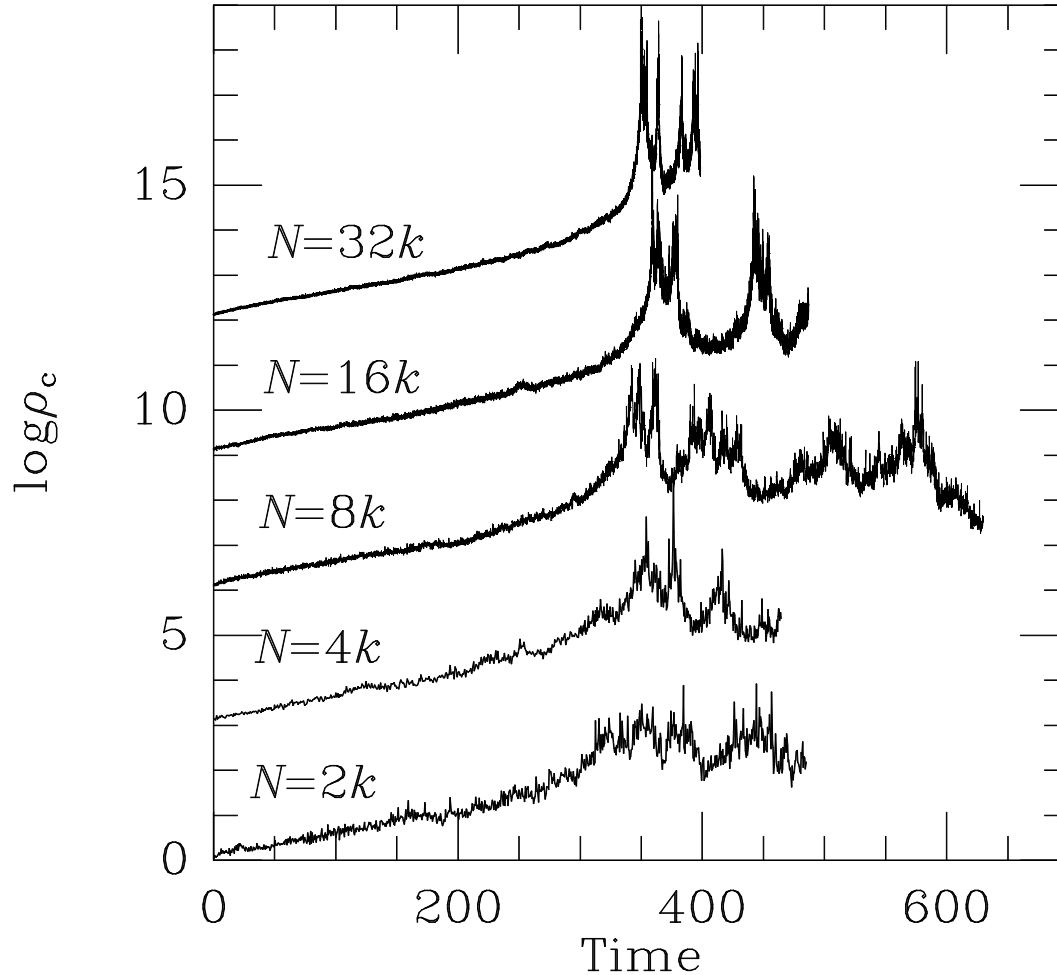
Sugimoto and
Bettwieser (1983)



Three curves:
Different energy
production
coefficients
(Different N)

N -body simulation

1996 : Confirmed with N -body simulation

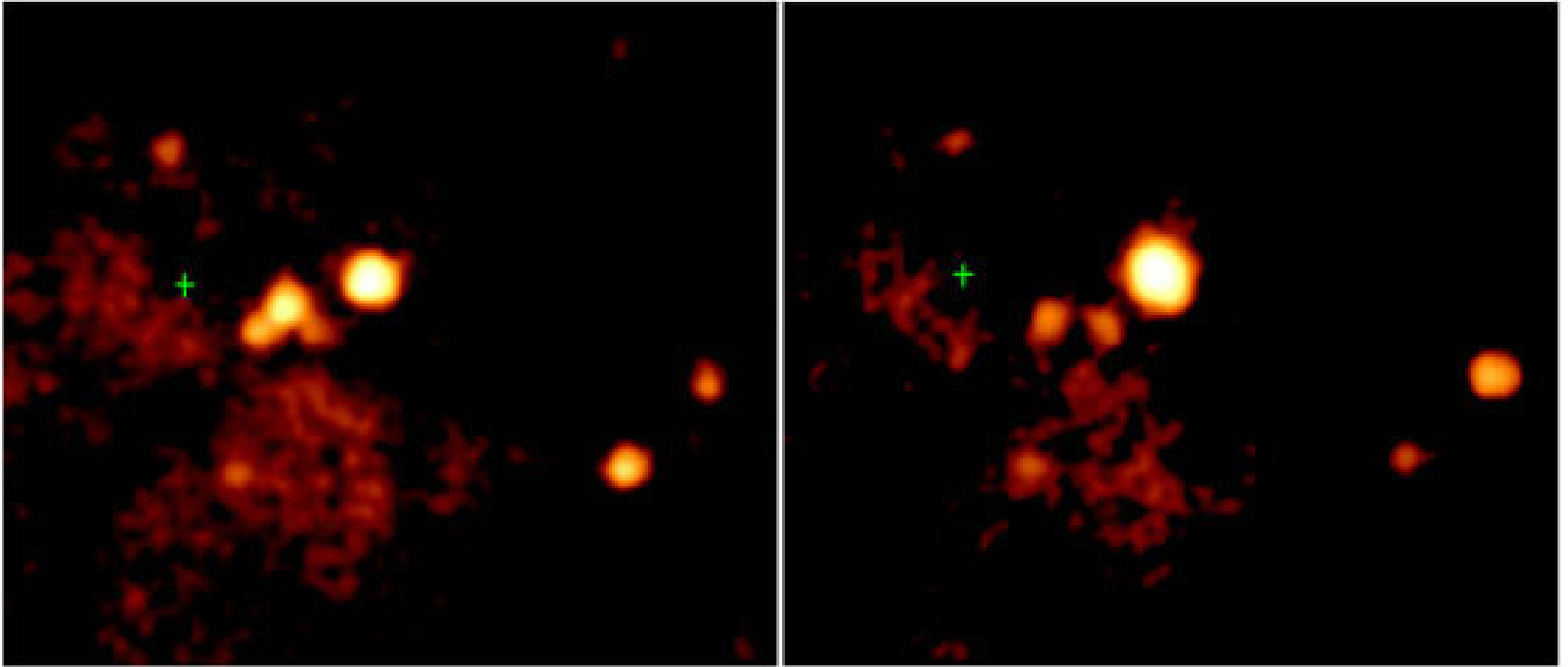


Some recent topics

- Collision of stars and blackhole formation

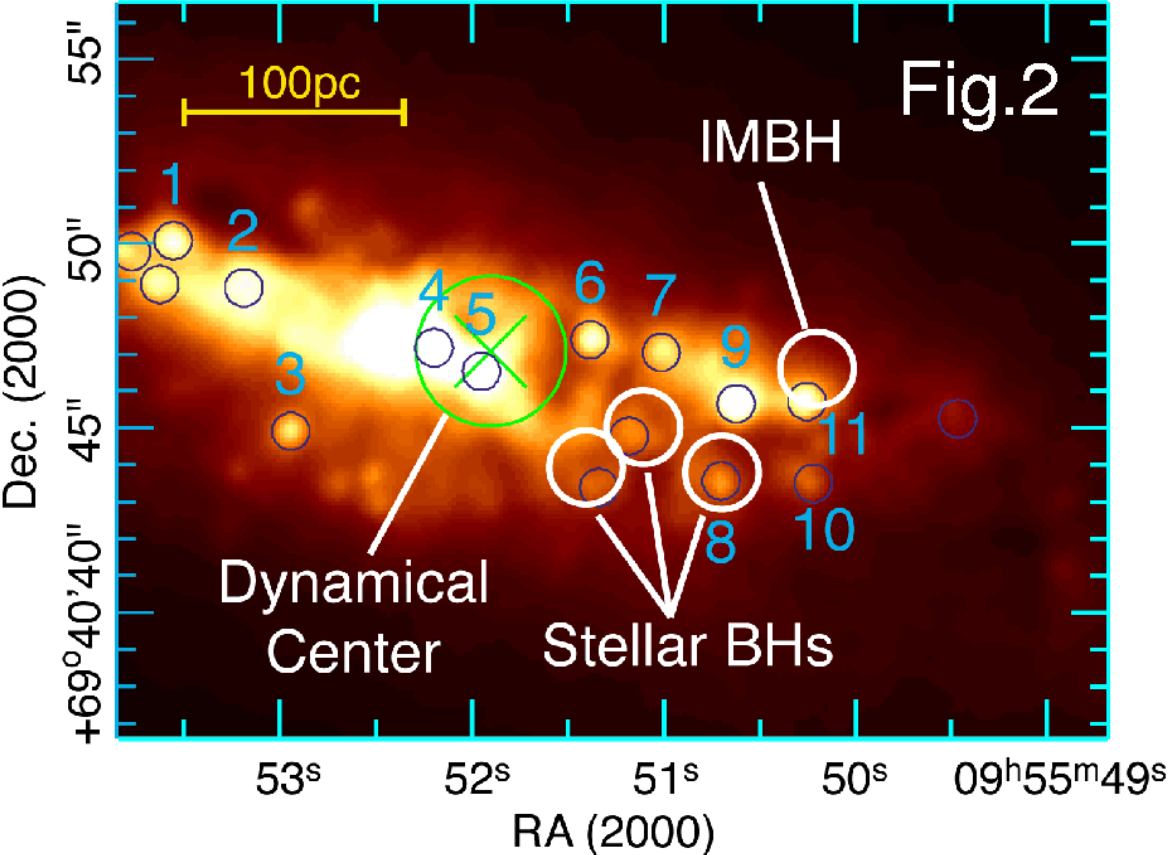
IMBH candidate in M82

Matsumoto *et al.* ApJL 547, L25



X-ray sources. Brightest one: Corresponds to 1000-solar-mass black hole (Intermediate-mass BH)

Subaru observation



Host star cluster

McCrady et al. 2003 (astro-ph/0306373)

Cluster #11 (MGG-11)

- $\sigma_r = 11.4 \pm 0.8 \text{ km/s}$
- half-light radius $1.2 \pm 0.17 \text{ pc}$
- kinetic mass $3.5 \pm 0.7 \times 10^5 M_\odot$
- Age $\sim 10 \text{ Myrs.}$

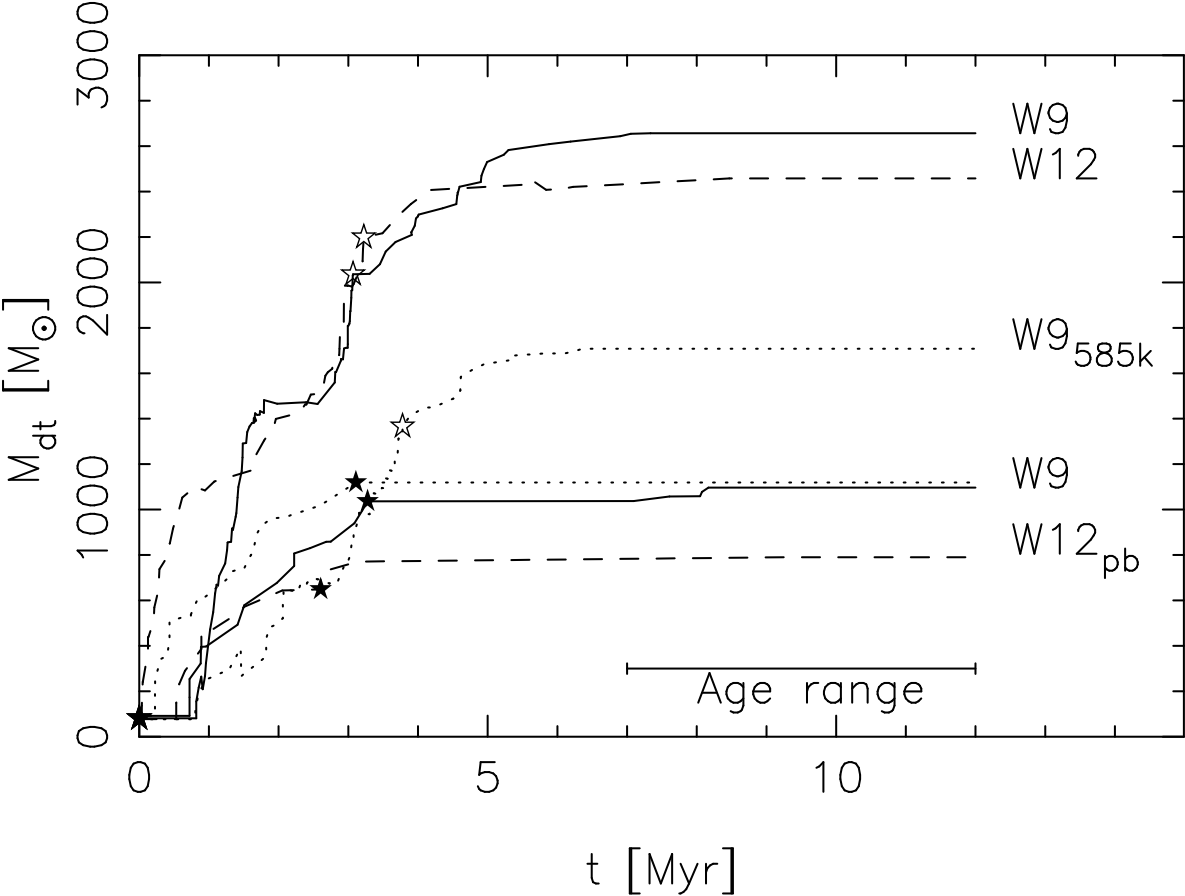
Very short relaxation time (less than 10Myrs)

A possible scenario

1. Massive stars sink to the center through relaxation
2. Supermassive star forms through collision and merging
3. This star collapses to a BH

Simulation

Portegies Zwart et al., Nature, Apr 15, 2004

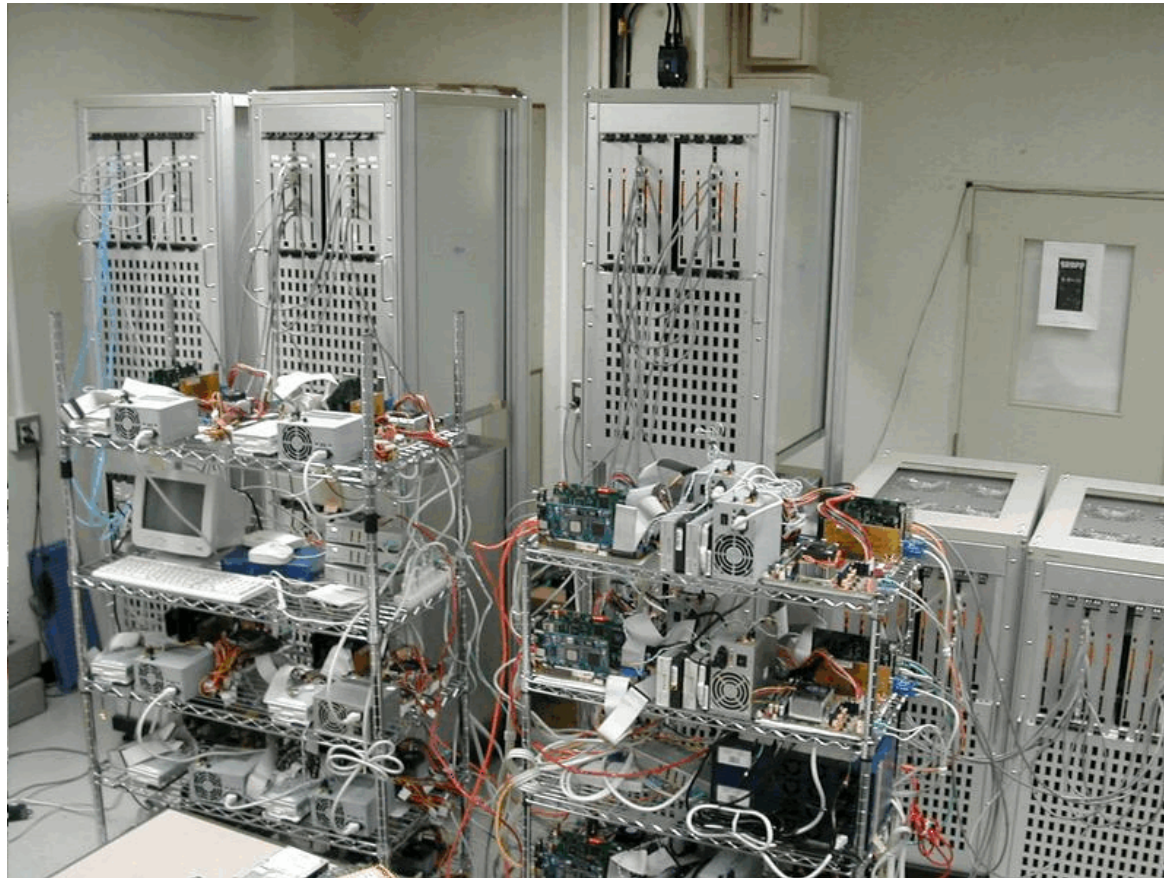


This scinario seems to work

GRAPE-6

Special-purpose computer for N -body simulations.

64 Tflops peak — World's **fastest** computer as of 2002.

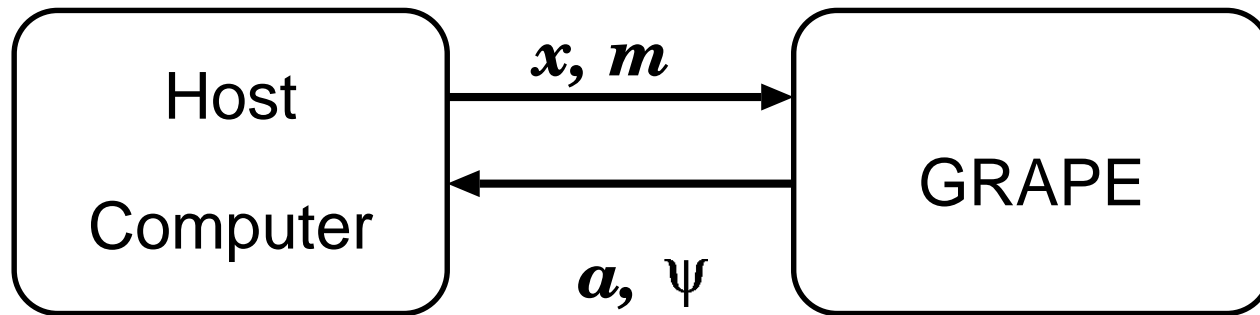


Basic idea of GRAPE

Special-purpose hardware for force calculation

General-purpose host for all other calculation

$$\frac{d^2 \mathbf{x}_i}{dt^2} = \sum_{j \neq i} -Gm_j \frac{\mathbf{x}_i - \mathbf{x}_j}{|\mathbf{x}_i - \mathbf{x}_j|^3}$$



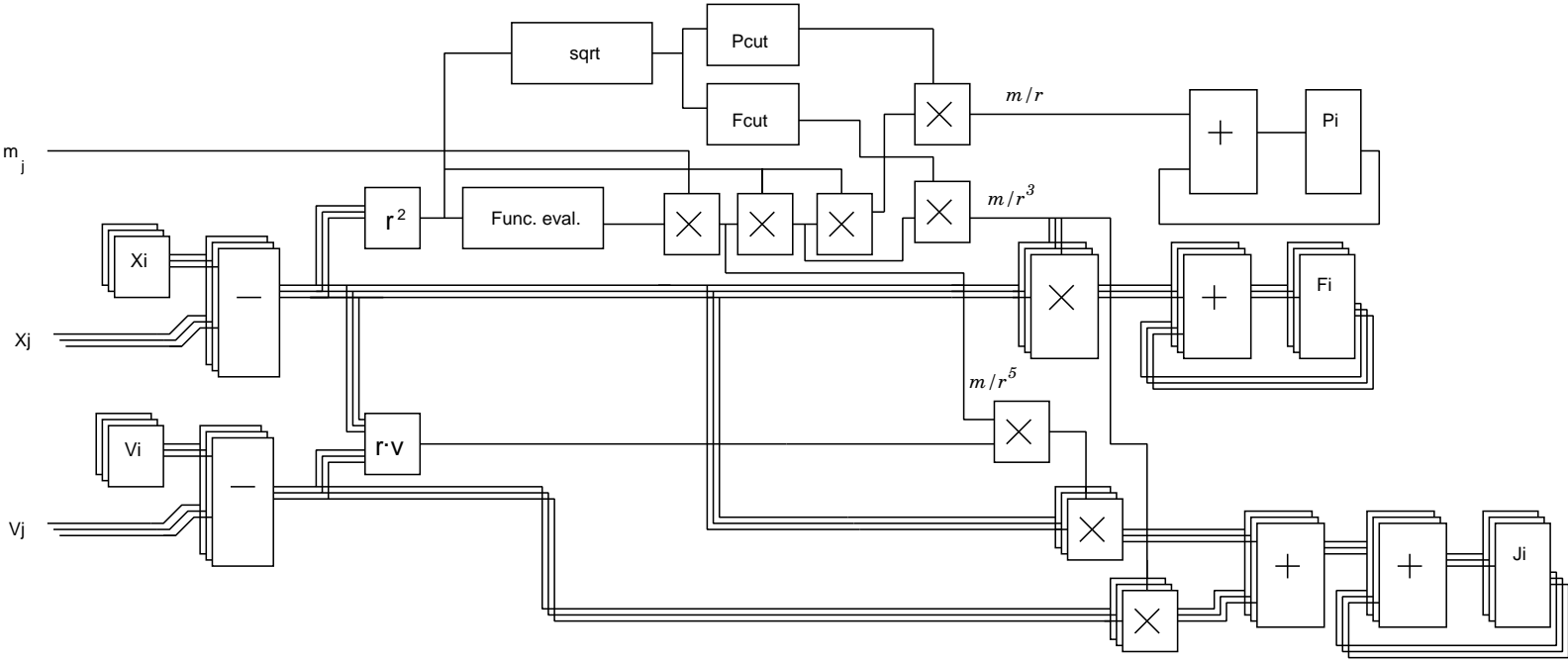
Time integration,
IO, etc

Flexibility

Force calculation

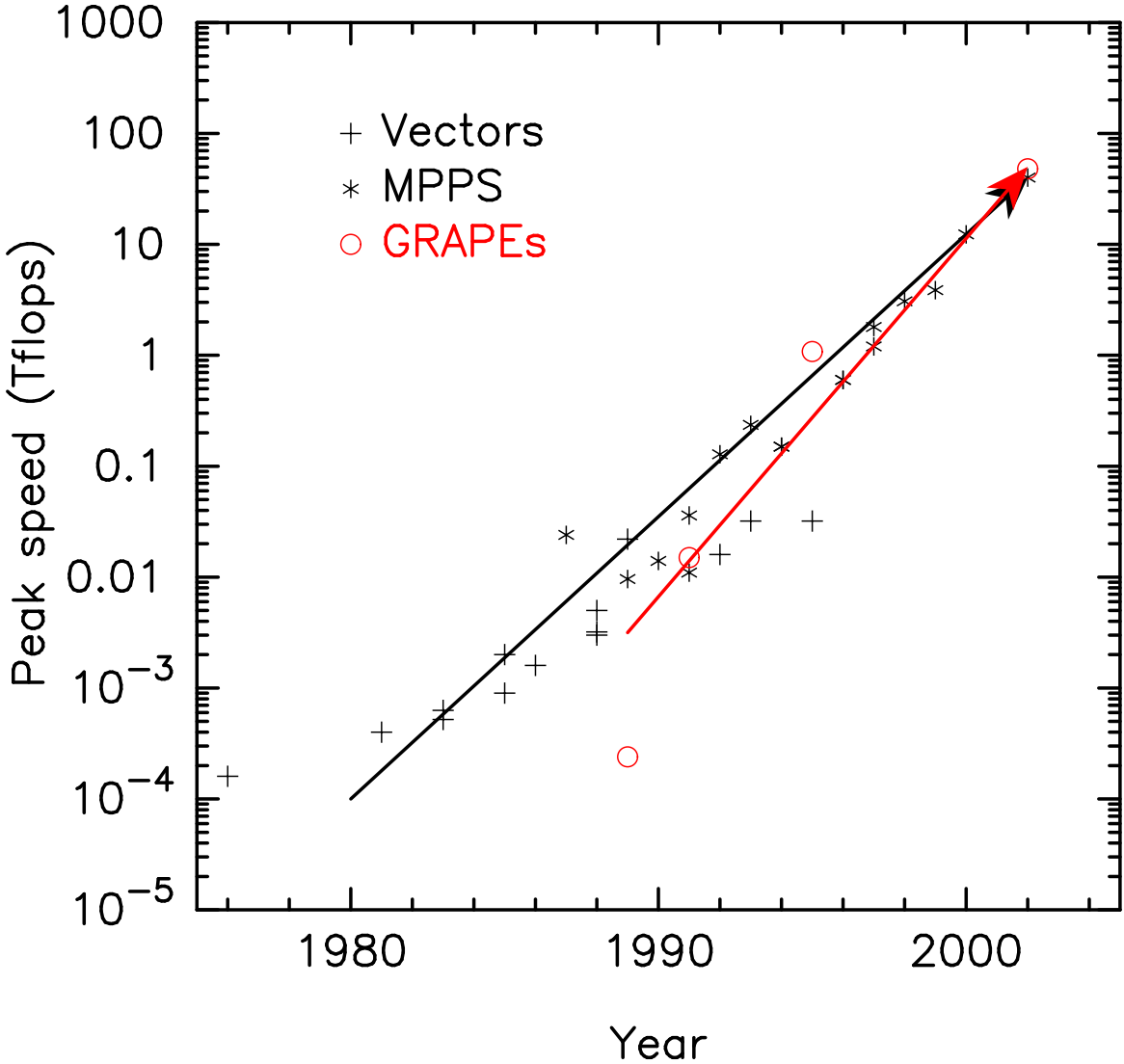
High performance

GRAPE-6 Processor pipeline



Calculates gravitational force, its first time derivative and potential.

Evolution of GRAPE systems

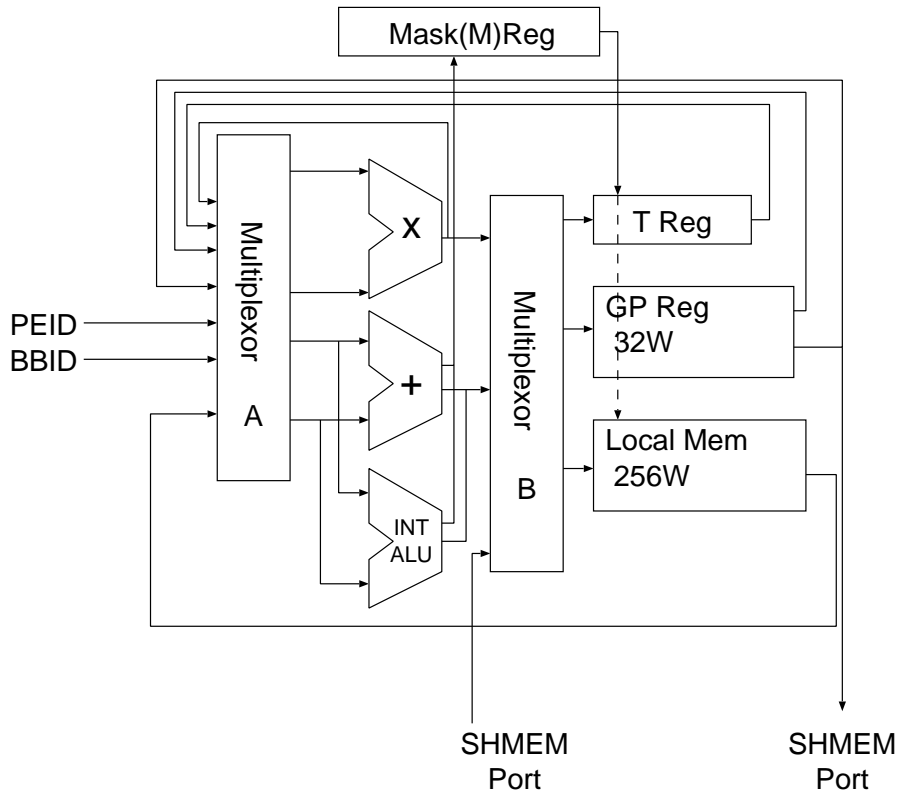


Next-Generation GRAPE

— GRAPE-DR

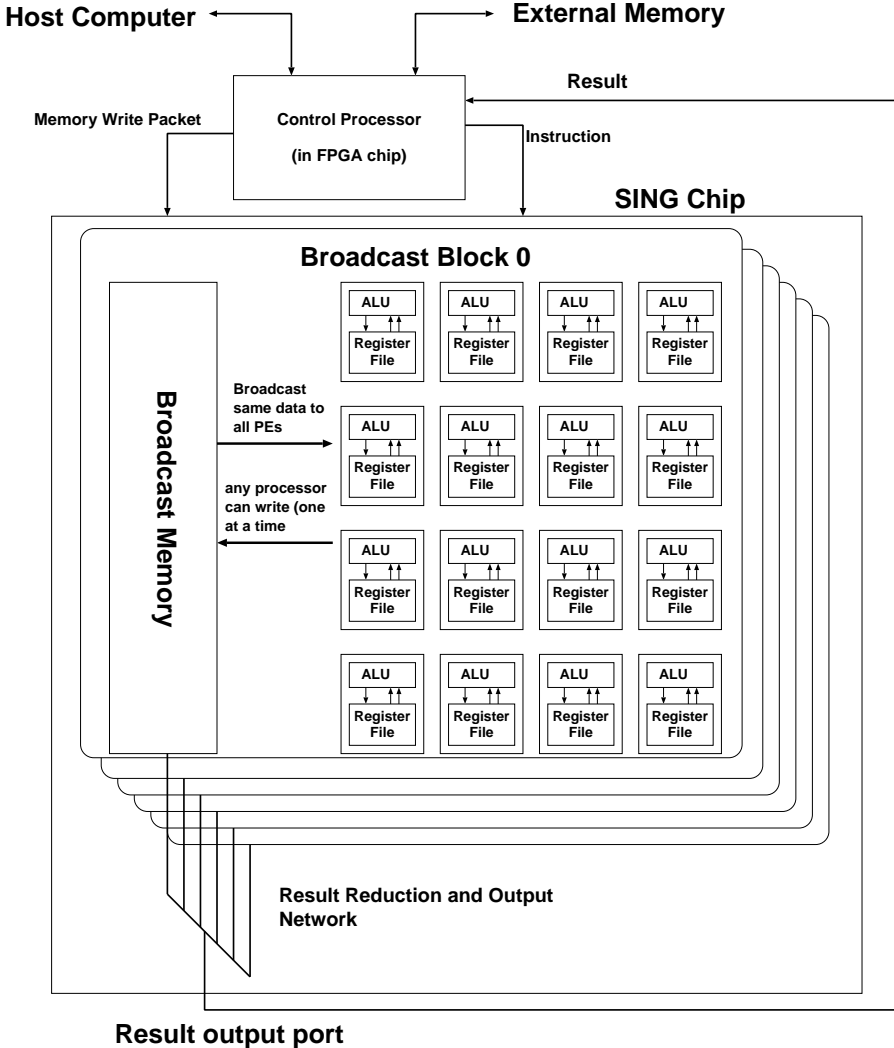
- Planned peak speed: 2 Pflops
- **New architecture — wider application range than previous GRAPEs**
- primarily to get funded
- No force pipeline. SIMD programmable processor
- Planned completion year: FY 2008 (early 2009)

Processor architecture



- Float Mult
- Float add/sub
- Integer ALU
- 32-word registers
- 256-word memory
- communication port

Chip structure



Collection of small processors.

512 processors on one chip
500MHz clock

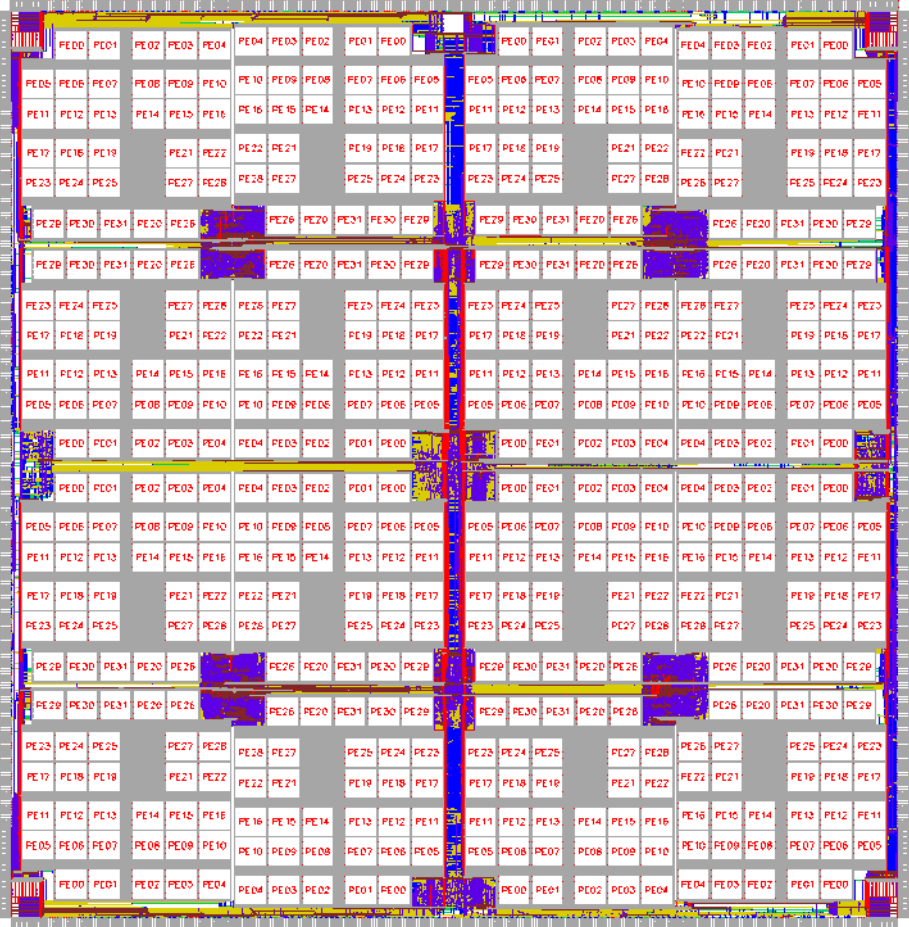
Peak speed of one chip: **0.5 Tflops** (20 times faster than GRAPE-6).

Development status



Sample chip delivered May 2006

Chip layout



- 32PEs in 16 groups
- 18mm by 18mm

Prototype board



2nd prototype board. (Designed by Toshi Fukushima)

Difference from the 1st one:

PCI-Express x8 interface

On-board DRAM

Designed to run real applications

Summary

- Gravitational Many-body problem is fun
- Recently, numerical integration of N -body systems have become useful tool for modelling observations.
- Special-purpose computers help.